

# Formal Proof: The Torsional Trefoil Markov Pump Theorem (TTMPT)

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June 19, 2026

## Abstract

The Torsional Trefoil Markov Pump Theorem (TTMPT) establishes a deterministic mathematical framework for a one-way stochastic pump mechanism acting on a longitudinal binding-scalar field. This paper formalizes the 10-step macro-cycle operating over a 12-channel lattice Hilbert space ( $\mathcal{H}_{\Lambda_{12}}$ ). By introducing dynamic transverse variables ( $\pi^2$  and  $\sqrt{\pi}$ ) as orthogonal heat sinks, formally postulating the projection of Hilbert-space amplitudes into classical Markov probabilities, and utilizing a topological “Tesla-valve” Markovian kernel, we rigorously prove that the system breaks detailed balance. Utilizing explicit stochastic boundary conditions, we demonstrate that the TTMPT generates a strictly positive, continuous net-forward probability current, advancing the binding scalar field while resetting the modular rotational phase to identity.

## 1 Axiomatic Foundations and State Space

### 1.1 The Rigorous State Space

The structural bedrock of the system consists of a triple-trefoil skeleton. We rigorously define the localized state space as a 12-channel lattice Hilbert space:

$$\mathcal{H}_{\Lambda_{12}} \cong \mathbb{C}^{12}.$$

The space operates using an explicit basis representing three coupled 4-phase streams (denoted by phase groups  $\alpha, \beta, \gamma$ ). A state vector at macro-step  $n$  is defined as the superposition

$$|\psi_n\rangle = \sum_{g \in \{\alpha, \beta, \gamma\}} \sum_{r=1}^4 c_{g,r}^{(n)} |e_{g,r}\rangle.$$

This state vector is subject to the standard normalization constraint ensuring conservation of total probability amplitude:

$$\sum_{g,r} |c_{g,r}^{(n)}|^2 = 1.$$

### 1.2 Dynamical Transverse Variables

To govern the continuous spatial scaling and continuous wave energy orthogonal to the central longitudinal binding field  $B$ , we introduce a scalar deformation parameter  $\delta$  alongside continuous

topological perturbation modes  $\hat{u}(t)$  and  $\hat{v}(t)$ . The transverse orthogonal foliation variables are rigorously defined as

$$\begin{aligned} u(t) &= \pi^2 + \delta \hat{u}(t) && \text{(Outer Manifold Variance),} \\ v(t) &= \sqrt{\pi} + \delta \hat{v}(t) && \text{(Inner Skeleton Winding Parameter).} \end{aligned}$$

These continuously parameterized variables dictate the continuous deformation of the Clifford torus during the stochastic pump cycle. They explicitly define the structural topological resistance function  $A_{ij}(t) = A_{ij}(u(t), v(t)) > 0$  for each graph edge, acting as independent orthogonal heat sinks for reverse-propagating current.

## 2 Topological Asymmetry and Broken Detailed Balance

### 2.1 The Postulated Hilbert-to-Markov Probability Bridge

To connect the quantum-like state vectors of  $\mathcal{H}_{\Lambda_{12}}$  to classical stochastic dynamics, we explicitly define the projection mapping.

**Definition** (Classical Markov Projection). The *classical Markov projection* of the Hilbert-space state is the probability vector  $p^{(n)}$  with entries

$$p_{g,r}^{(n)} = |c_{g,r}^{(n)}|^2.$$

Consolidating these probabilities yields a 12-dimensional classical probability vector  $p^{(n)} = (p_1^{(n)}, \dots, p_{12}^{(n)})$ . We postulate that, under the coarse-grained pump dynamics, this projected probability vector evolves according to a row-stochastic Markov kernel  $K$  such that

$$p^{(n+1)} = p^{(n)} K.$$

### 2.2 Broken Detailed Balance and the Asymmetric Kernel

For a finite Markov chain with stationary distribution  $\pi$ , *detailed balance* holds if  $\pi_i K_{ij} = \pi_j K_{ji}$  for all states  $i, j$ . To break detailed balance, we define a non-empty oriented forward edge set  $E^+ \neq \emptyset$ .

**Lemma 2.1** (Asymmetric Markovian Kernel). *Suppose that for every  $(i, j) \in E^+$  the row-stochastic transition probabilities satisfy*

$$K_{ij} > 0, \quad K_{ji} = K_{ij} e^{-\lambda A_{ij}},$$

where  $\lambda > 0$  and  $A_{ij} > 0$ . If, additionally, the stationary distribution satisfies the structural condition

$$\pi_i > \pi_j e^{-\lambda A_{ij}},$$

then it strictly follows that

$$\pi_i K_{ij} - \pi_j K_{ji} > 0.$$

Hence the net probability current  $J = \sum_{(i,j) \in E^+} (\pi_i K_{ij} - \pi_j K_{ji})$  over the non-empty set is strictly positive.

*Proof.* The current along any directed edge  $(i, j)$  is  $J_{ij} = \pi_i K_{ij} - \pi_j K_{ji}$ . Substituting the topological suppression rule  $K_{ji} = K_{ij} e^{-\lambda A_{ij}}$  gives

$$J_{ij} = K_{ij} (\pi_i - \pi_j e^{-\lambda A_{ij}}).$$

Because we have explicitly assumed  $K_{ij} > 0$  and  $\pi_i > \pi_j e^{-\lambda A_{ij}}$ , the term in parentheses is strictly positive. Therefore  $J_{ij} > 0$  for every edge in  $E^+$ . Since  $E^+ \neq \emptyset$ , there exists at least one edge  $(i, j)$  such that  $\pi_i K_{ij} \neq \pi_j K_{ji}$ , so detailed balance fails. Consequently the net current  $J$  over the non-empty set is a sum of strictly positive terms.  $\square$

### 3 Mechanics of the 10-Step Macro-Cycle

#### 3.1 Phase I: Compression and the Frame 5 Solver

Let the 12-channel state lie in a Hilbert space  $\mathcal{H} \cong \mathbb{C}^{12}$ , and assume  $\|\psi_0\rangle\| = 1$ . Let  $P_1, \dots, P_4$  be orthogonal projections on  $\mathcal{H}$  and let  $\Pi$  be a unitary discrete phase operator. Define the compression operator

$$C = P_4 P_3 \Pi P_2 P_1.$$

Since each  $P_k$  is norm non-increasing and  $\Pi$  is norm-preserving,  $C$  is norm non-increasing ( $\|C|\psi\rangle\| \leq \| |\psi\rangle \|$ ). Assuming  $C|\psi_0\rangle \neq 0$ , define the normalized pre-pump compressed state by

$$|\psi_{\text{comp}}\rangle = \frac{C|\psi_0\rangle}{\|C|\psi_0\rangle\|} = \frac{P_4 P_3 \Pi P_2 P_1 |\psi_0\rangle}{\|P_4 P_3 \Pi P_2 P_1 |\psi_0\rangle\|}.$$

Then  $\| |\psi_{\text{comp}}\rangle \| = 1$ . The Frame 5 Algebraic Solver provides the geometric phase-lock parameter

$$s = 1 - \frac{(Q_\alpha + Q_\beta - Q_\gamma)^2}{4 Q_\alpha Q_\beta}, \quad Q_\alpha, Q_\beta > 0.$$

Setting  $s = \sin^2\left(\frac{2\pi}{3}\right) = \frac{3}{4}$  imposes the symmetric  $120^\circ$  phase-lock condition. This geometric condition motivates the topology of the cycle, while the rigorous contraction mechanism is supplied by the Markov kernel assumptions in the pump phase.

#### 3.2 Phase II: Core Pump Stroke and Time Parameterization

During the pump interval  $t \in [t_5, t_7]$ , let  $K(t)$  be a continuously parameterized family of stochastic transition matrices. For example, in the row-stochastic convention,  $K_{ij}(t) \geq 0$  and  $\sum_j K_{ij}(t) = 1$ . The compressed state induces a probability vector

$$p_i^{\text{comp}} = |\langle e_i | \psi_{\text{comp}} \rangle|^2.$$

The Markov kernel evolves this probability layer during the pump stroke. If the larger theorem requires strict contraction, assume there exists  $0 \leq \lambda < 1$  such that  $\|pK(t) - qK(t)\|_1 \leq \lambda \|p - q\|_1$  for all admissible probability vectors  $p, q$  and all  $t \in [t_5, t_7]$ . Let  $S_1, S_2, S_3$  be unitary rotational operators ( $S_k^* S_k = I$ , so  $S_k^{-1} = S_k^*$ ). Assuming  $\sum_{k=1}^3 (S_k + S_k^{-1})|\psi_0\rangle \neq 0$ , define the normalized quaternion centroid by

$$|\Psi_4\rangle = \frac{\sum_{k=1}^3 (S_k + S_k^{-1})|\psi_0\rangle}{\|\sum_{k=1}^3 (S_k + S_k^{-1})|\psi_0\rangle\|}.$$

The active current  $J(t)$  is then interpreted as a parameterized flow directed toward  $|\Psi_4\rangle$ .

#### 3.3 Phase III: Double Toroidal Collapse

Let  $(u(t), v(t)) \in \mathbb{T}^2$  parameterize the toroidal collapse, and let  $A_{ij}: \mathbb{T}^2 \rightarrow \mathbb{R}_{\geq 0}$  be continuous topological resistance functions. Introduce a reverse-energy functional  $E_{\text{rev}}(t) \geq 0$ . The collapse phase is dissipative if

$$\frac{d}{dt} E_{\text{rev}}(t) = - \sum_{i,j} A_{ij}(u(t), v(t)) R_{ij}(t)^2 \leq 0.$$

It is strictly dissipative whenever at least one active resistance mode satisfies  $A_{ij}(u(t), v(t)) R_{ij}(t)^2 > 0$ . Thus reverse variance is dissipated into transverse modes while the longitudinal field remains protected by a separately imposed conservation or lower-bound condition.

### 3.4 Phase IV: Encapsulation and Phase Reset

Let  $\Theta_k \in \mathbb{R}/2\pi\mathbb{Z}$  and define the modular phase update by  $\Theta_{k+1} = \Theta_k + \omega_k$ . After one full 10-step macro-cycle,

$$\Theta_{10} = \Theta_0 + \sum_{k=0}^9 \omega_k.$$

Encapsulation requires the residual modular rotation to vanish:  $\sum_{k=0}^9 \omega_k \equiv 0 \pmod{2\pi}$ . Equivalently,  $\Theta_{10} \equiv \Theta_0 \pmod{2\pi}$ .

## 4 The Torsional Trefoil Markov Pump Theorem (TTMPT)

With the axiomatic mechanics established, we formalize the principal theorem.

**Theorem 4.1 (TTMPT).** *Let  $\Lambda_{12}$  be a finite state space with 12 states, and let  $E^+ \subseteq \Lambda_{12} \times \Lambda_{12}$  be a non-empty oriented forward edge set ( $E^+ \neq \emptyset$ ). Let  $[t_5, t_7]$  be a non-degenerate compact interval with  $t_7 > t_5$ . For each  $t \in [t_5, t_7]$ , let  $K(t)$  be a continuously parameterized family of discrete-time row-stochastic, irreducible, aperiodic transition matrices. Assume that  $K_{ij}(t)$ , the unique stationary distribution  $\pi_i(t)$ , and the topological resistance function  $A_{ij}(t) = A_{ij}(u(t), v(t))$  are continuous in  $t$ , satisfying*

$$\pi(t)K(t) = \pi(t), \quad \sum_i \pi_i(t) = 1, \quad \pi_i(t) > 0.$$

Assume that for every  $(i, j) \in E^+$  and every  $t \in [t_5, t_7]$ ,

$$K_{ij}(t) > 0, \quad K_{ji}(t) = K_{ij}(t) e^{-\lambda A_{ij}(t)},$$

where  $\lambda > 0$  and  $A_{ij}(t) > 0$ , and assume

$$\pi_i(t) > \pi_j(t) e^{-\lambda A_{ij}(t)}.$$

Define the continuous forward probability current by

$$J(t) = \sum_{(i,j) \in E^+} (\pi_i(t)K_{ij}(t) - \pi_j(t)K_{ji}(t)).$$

Then  $J(t) > 0$  for all  $t \in [t_5, t_7]$ . Consequently, if  $B$  satisfies  $\frac{dB}{dt} = J(t)$ , then

$$B(t_7) > B(t_5).$$

Furthermore, if the discrete phase update is  $\Theta_{k+1} = \Theta_k + \omega_k$  and satisfies the closure condition  $\sum_{k=0}^9 \omega_k \equiv 0 \pmod{2\pi}$ , then

$$\Theta_{10} \equiv \Theta_0 \pmod{2\pi}.$$

*Proof.* For each  $(i, j) \in E^+$  at any time  $t \in [t_5, t_7]$ , using the topological suppression rule,

$$\begin{aligned} \pi_i(t)K_{ij}(t) - \pi_j(t)K_{ji}(t) &= \pi_i(t)K_{ij}(t) - \pi_j(t)K_{ij}(t) e^{-\lambda A_{ij}(t)} \\ &= K_{ij}(t) (\pi_i(t) - \pi_j(t) e^{-\lambda A_{ij}(t)}). \end{aligned}$$

Since  $\pi_i(t) > \pi_j(t) e^{-\lambda A_{ij}(t)}$  by hypothesis, the quantity in parentheses is strictly positive, and factoring out the strictly positive  $K_{ij}(t)$  yields

$$\pi_i(t)K_{ij}(t) - \pi_j(t)K_{ji}(t) > 0.$$

Because  $E^+$  is non-empty ( $E^+ \neq \emptyset$ ), there exists at least one edge  $(i, j)$  such that  $\pi_i(t)K_{ij}(t) \neq \pi_j(t)K_{ji}(t)$ , so detailed balance fails at each  $t \in [t_5, t_7]$ . Summing over the forward edge set gives a strictly positive continuous net current:

$$J(t) = \sum_{(i,j) \in E^+} (\pi_i(t)K_{ij}(t) - \pi_j(t)K_{ji}(t)) > 0.$$

Given that the scalar field evolves as  $\frac{dB}{dt} = J(t)$ , we integrate over the non-degenerate interval  $[t_5, t_7]$ :

$$B(t_7) - B(t_5) = \int_{t_5}^{t_7} J(t) dt.$$

Because  $J(t) > 0$  and is continuous throughout the entire interval, and  $t_7 > t_5$ , the definite integral is strictly positive:

$$\int_{t_5}^{t_7} J(t) dt > 0.$$

Hence  $B(t_7) - B(t_5) > 0$ , demonstrating  $B(t_7) > B(t_5)$ .

Finally, iterating the discrete phase updates through the 10-step macro-cycle gives

$$\Theta_{10} = \Theta_0 + \sum_{k=0}^9 \omega_k.$$

Using the phase-closure constraint  $\sum_{k=0}^9 \omega_k \equiv 0 \pmod{2\pi}$ , we obtain

$$\Theta_{10} \equiv \Theta_0 \pmod{2\pi}.$$

Therefore, one full macro-cycle produces strictly positive advancement of the scalar field  $B$  while successfully resetting the modular phase to identity.  $\square$

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*End of Proof.*